

Studies of track finding for long-lived particles at STCF*

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Reconstruction of the trajectories of charged particles at High Energy Physics experiments is a complicated task, in particular those of long-lived particles. At the future Super Tau-Charm Facility (STCF), long-lived particles are present in several important benchmark physics processes. A Common Tracking Software was used to reconstruct the trajectories of long-lived particles and it is shown that the track finding performance of the commonly used Combinatorial Kalman Filter for long-lived particles is limited by the seeding algorithm. This can be improved by steering the Combinatorial Kalman Filter with initial tracks provided by Hough Transform. The track finding performance of combined Hough Transform and Combinatorial Kalman Filter evaluated using the process $J/\psi \rightarrow \Lambda(\rightarrow p\pi^-)\bar{\Lambda}(\rightarrow \bar{p}\pi^+)$ at STCF is presented.

Keywords: Track finding, Common tracking software, Hough Transform, Long-lived particles

I. INTRODUCTION

Standard Model (SM) [1, 2] of particle physics including the unified Electro-Weak (EW) and Quantum Chromodynamics (QCD) theories, has explained successfully almost all experimental results about the microscopic world. However, a couple of questions still remain, e.g. baryon asymmetry of the universe, dark matter, neutrino masses, number of flavors. Beijing Electron Positron Collider (BEPCII) - Beijing Spectrometer (BESIII) [3] is the only multi-GeV e^+e^- collider operating in the τ -charm sector, which provides an unique platform for studying non-perturbative QCD and strong interactions of the SM. The Super Tau-Charm Facility (STCF) [4, 5] is designed to continue and extend the physics programs at BEPCII in near future, including probing the nature of the strong interactions and hadron structure, precise inspection of electroweak theories, exploring the asymmetry of matter-antimatter and searching for new physics beyond the SM. STCF will operate at a center-of-mass-energy of 2-7 GeV and a peak luminosity above $0.5 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$ [6], which is two orders higher than that at BEPCII.

The reconstruction of charged particles is the most fundamental and critical step in the data processing chain of High Energy Physics (HEP) experiments. To fulfill the physics goals and to further maximize the physics potential at the STCF, the charged particles need to be reconstructed with good efficiency. This includes not only those particles that decay immediately upon production but also the long-lived particles [7], e.g. the Λ and Ξ hyperons, which are relevant with a couple of important physics goals at STCF. For example, the weak decays of the Λ and Ξ hyperons provide promising channels for searching for new sources of CP violation [8–10]. The hyperon samples at STCF can also be

used to measure the time-like nucleon and hyperon form factors for Q^2 values as high as 40 GeV 2 [5]. Meanwhile, it is quite challenging to reconstruct the trajectories of long-lived particle decay products because the long-lived particles may decay within or outside the inner tracker hence having very limited number of hits recorded by the inner tracker.

The Kalman Filter (KF) [11] algorithm is the most commonly used algorithm for tracking in HEP and nuclear physics [12]. The Combinatorial Kalman Filter (CKF) [13, 14] is an extended version of the KF, where the measurements are progressively added to the track during the track propagation steered by an initial estimate of the track parameters, i.e. seed. The impact of magnetic field and material effects is incorporated during track propagation hence CKF is capable to resolve the hit ambiguity in a very dense tracking environment. For this reason, CKF is deployed to find tracks by several experiments e.g. ATLAS [15] and CMS [16], where thousands of tracks are present in a single event. CKF is also the primary track finding algorithm at BelleII experiment [17]. Recently, the CKF algorithm developed by BelleII experiment was reused to study the tracking performance [18] at the Circular Electron-Positron Collider (CEPC) [19]. Despite the great advantages of CKF, one downside of the KF-based tracking algorithms is that they are subject to the performance of the seeding algorithm, which might provide poor performance for long-lived particles. Recently, a track finding algorithm based on the Hough Transform used by BelleII experiment [20] and BESIII experiment [21] has been developed at STCF [22], where tracking performance for prompt particles without vertex displacement have been studied. It demonstrates promising tracking performance, in particular good robustness against local hit inefficiency. However, the tracking efficiency can be deteriorated by the presence of background hits at low transverse momentum.

The ACTS (A Common Tracking Software) [23, 24] is an emerging open-source tracking software for HEP and nuclear physics experiments, with a suite of detector-agnostic and framework-independent modular track and vertex reconstruc-

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tion algorithms. The promising performance of the KF and CKF algorithms in ACTS is underscored by their widespread adoption by experiments such as FASER [25], sPHENIX [26] and a few R&D studies at STCF [27] and BESIII [28]. Notably, ACTS has demonstrated its generality across a series of tracking detector types [29]. However, the performance of ACTS for long-lived particles has not been investigated.

In this study, the tracking performance of STCF with a fully gaseous tracking system consisting of a μ -RWELL [30]-based inner tracker and a drift chamber using combined Hough Transform and ACTS CKF to boost the tracking performance for long-lived particles is studied. The manuscript is organized as follows. Section II presents a brief introduction of the STCF detector. In Section III, the tracking workflow with different algorithms is introduced. Section IV focuses on tracking performance for benchmark process with long-lived particles at STCF. A brief conclusion is given in Section V.

II. STCF DETECTOR

The STCF detector [5] ensures comprehensive coverage of the solid angle encompassing the collision point, as depicted in Fig. 1. The STCF detector consists of a tracking system comprising an Inner Tracker (ITK) and a Main Drift Chamber (MDC), along with a Ring Imaging Cherenkov (RICH) detector [31] and a DIRC-like Time-of-Flight (DTOF) detector [32] for particle identification in the barrel and endcap regions. Additionally, it incorporates a uniform Electro-magnetic Calorimeter (EMC) [33], a superconducting solenoid magnet generating 1 Tesla axial magnetic field, and a Muon Detector (MUD) positioned at the detector system's outermost layer.

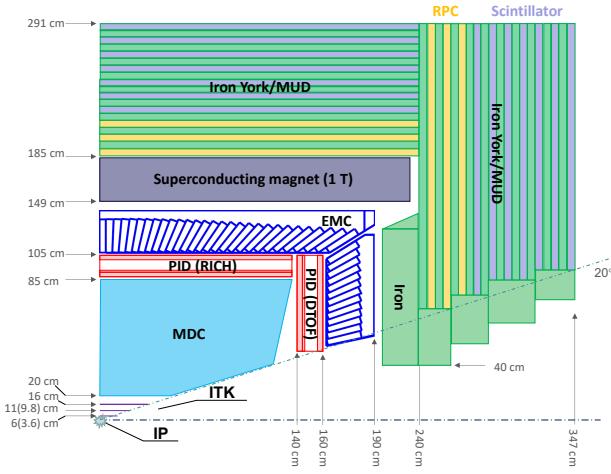


Fig. 1. Schematic layout of the STCF detector. The number in brackets indicate the radii of the MAPS-based ITK. Figure is taken from Ref. [5].

To ensure optimal tracking efficiency for low-momentum charged particles, the ITK within the tracking system cover-

ing a polar angle range of 20° to 160° (i.e. $l \cos\theta < 0.94$) comprises three layers of low-material budget silicon or gaseous detectors using either MAPS-based or μ -RWELL-based technology [30]. This study focuses on the μ -RWELL-based ITK with the three layers placed at an inner radii of 60 mm, 110 mm, and 160 mm, respectively, and each layer has a thickness of approximately 6.5 mm. It provides a spatial resolution around $100 \mu\text{m}$ in the r - ϕ direction and around $400 \mu\text{m}$ in the z direction. For the MAPS-based ITK, the radii of the three silicon layers is 36 mm, 98 mm and 160 mm, respectively, and a hit resolution of $30 \mu\text{m} \times 180 \mu\text{m}$ is assumed. Unless explicitly specified, ITK denotes the μ -RWELL-based ITK by default.

Central to the tracking system of the STCF detector, the Main Drift Chamber (MDC) operates using a He/C₃H₈(60/40) gas mixture and features a square cell configuration with a superlayer wire arrangement. The superlayers within the MDC alternate between stereo layers ("U" or "V") and axial layers, each containing six layers. In total, the MDC comprises eight superlayers (AUVAUVA) and 48 layers, with inner and outer radii of 200 mm and 850 mm, respectively. The MDC provides spatial resolutions ranging between $120 \mu\text{m}$ and $130 \mu\text{m}$.

III. TRACK RECONSTRUCTION USING COMBINED HOUGH TRANSFORM AND CKF

The workflow of track reconstruction using combined Hough Transform and ACTS CKF is illustrated in Fig. 2. The ACTS CKF is used to find track candidates through track fitting steered by the initial track parameters provided by either ACTS seeding algorithm or Hough Transform algorithm developed within the STCF offline software.

A. Interface between STCF offline software and ACTS

The Offline Software System of the Super Tau-Charm Facility (OSCAR) [34, 35] serves as the offline event processing framework for the STCF. It provides common services for data processing and a suite of application tools dedicated to event generation, simulation, reconstruction, and physics analysis. For simulation purposes, OSCAR incorporates generation of τ -charm physics processes facilitated by the KKMC [36] generator, while particle decays are modeled with EvtGen as used by BESIII experiment [37], both seamlessly integrated within the framework. The STCF detector geometry is described using the Detector Description Toolkit, DD4Hep [38], with all geometric parameters stored in compact files utilizing the eXtensible Markup Language (XML) [39]. To simulate the interaction of particles with the detector comprehensively, Geant4 [40] is integrated into OSCAR, ensuring a sophisticated full simulation. For track reconstruction, the track finding algorithm based on Hough Transform is developed in OSCAR.

The interface between OSCAR and ACTS facilitates the transforming of experimental geometry, measurements, and

156 initial track estimates into corresponding ACTS representations. Geometry plugins within ACTS are tailored to streamline the conversion of experimental geometry representations, such as DD4hep or TGeo [41], into ACTS internal geometry description. For ITK, the transformation involves converting the signal readout unit tube within each μ -RWELL layer into sensitive cylinder surfaces. Similarly, for MDC, the process entails transforming each sense wire within a drift cell into a line surface. Leveraging dedicated material mapping tools within ACTS, detailed material descriptions are projected onto internal auxiliary surfaces of the ACTS geometry. For the conversion of measurements and initial track parameters, two ROOT [42]-based readers have been developed. One reader extracts simulated hits from full simulation data and converts them into ACTS measurements taking into account the resolution of the detectors. Another reader converts the initial estimate of the track parameters provided by the Hough Transform algorithm into ACTS track parameters.

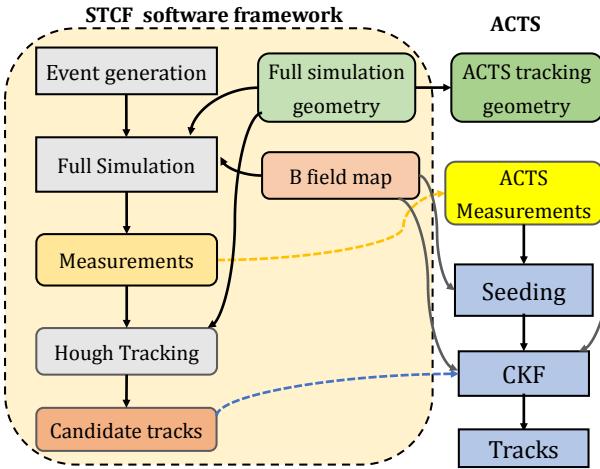


Fig. 2. The workflow of studying tracking performance using STCF software framework and ACTS.

B. ACTS seed finding

174 The seeding algorithm in ACTS aims to find a few measurements which can provide position coordinates (x, y, z) in 175 the global coordinate frame associated with a single particle 176 to initiate the track following process. Without a seed, a particle 177 cannot be reconstructed, hence the seed finding algorithm 178 aims to find at least one seed for each particle in the detector 179 acceptance region.

180 In a uniform magnetic field along the global z axis, the 181 helical trajectory of a charged particle is accurately defined 182 by three measurements, thus forming a seed. In the case of 183 STCF, these seeds are generated by combining three compatible 184 measurements from the ITK detector with one measurement 185 per ITK layer, as illustrated in Fig. 3. For each candidate 186 seed, the curvature and center of the circle on the x - y plane 187 are determined using the Conformal Transform [43, 44]. Sub- 188 189

sequently, these parameters are utilized to calculate the transverse momentum and the transverse impact parameter on the x - y plane, which are required to satisfy the criteria optimized for the relevant physics processes. The bending of the seed in the r - z plane is also required to be smaller than a threshold, which is optimized taking into account the impact of magnetic field and multiple scattering. See Ref. [45] for a detailed description of the ACTS seeding.

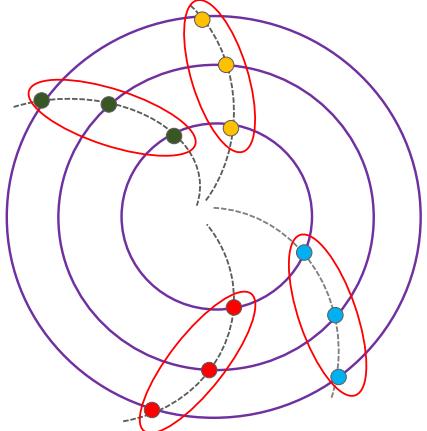


Fig. 3. Illustration of ACTS seeding using measurements from STCF ITK.

C. Track finding with Hough Transform

198 The principle of Hough Transform for track finding is illustrated in Fig. 4. With the presence of a magnetic field along 199 global z axis, the projection of the track in the geometrical 200 global x - y plane is a circle and the projection of the track 201 in the geometrical s - z plane (s is the path length of the track 202 in the x - y plane) is a straight line. The Conformal Trans- 203 form can convert the projection of a track in the transverse 204 x - y plane passing through the origin into a straight line, with 205 a drift circle tangent to the projection of the track converted 206 to another circle tangent to the straight line, in the Confor- 207 mal u - v space. For a displaced track which has non-zero but 208 relatively small transverse impact track parameter, i.e. d_0 , 209 compared to the radius of the circle projected in the geomet- 210 rical transverse x - y plane, as shown in Fig. 4 left panel, the 211 trajectory of the transformed measurements (either a point or 212 a drift circle) on the Conformal space can be approximated 213 parameterized using a straight line, as shown in Fig. 4 mid- 214 le panel. The Hough Transform for track finding operates 215 on the principle that a straight line in the geometrical or Con- 216 formal space can be described by two parameters, angle θ of 217 its normal and its algebraic distance ρ from the origin. For 218 a point with coordinate (u, v) on the straight line, it can be 219 transformed into a sinusoidal curve in the Hough space, i.e. 220 transformed into a sinusoidal curve in the Hough space, i.e. 221

$$\rho = u \cdot \cos\alpha + v \cdot \sin\alpha \quad (1)$$

222 , and for a circle with center (u, v) and radius r tangent to the

224 straight line, it can be transformed into two sinusoidal curves
225 in the Hough space, i.e.

$$226 \quad \rho = u \cdot \cos\alpha + v \cdot \sin\alpha \pm r \quad (2)$$

227 The process of finding the measurements or drift circles
228 that arise from the same track in either the Conformal u - v
229 space or the geometrical s - z space becomes identifying the
230 curves that have an intersection in Hough space as shown in
231 Fig. 4 right panel, i.e. a 2D histogram with optimized binning
232 taking into account the resolution of the track parameters are
233 filled if a curve passes through it and the peaks of the his-
234 tograms are identified, as detailed in Ref. [22]. The param-
235 eters at the intersection can be converted to the track param-
236 eters describing the track projected to either the geometrical
237 transverse x - y plane or the geometrical s - z plane.

238 The workflow of track finding using Hough Transform in
239 OCSCAR is shown in Fig. 5. Initially, the measurements from
240 ITK and MDC axial wires are used to reconstruct the projec-
241 tions of the tracks on the x - y plane, denoted as 2D tracks,
242 followed by circle fitting to extract track parameters of the
243 2D tracks. This is succeeded by associating the MDC stereo
244 wire measurement candidates to the 2D tracks, where the z
245 position and path length s of the track at the stereo wires are
246 derived simultaneously. For each stereo wire measurement,
247 two z position solutions can be obtained, and measurements
248 from other tracks may be wrongly assigned to a 2D track.
249 Therefore, a secondary application of Hough Transform is
250 employed to find the tracks in the s - z plane. More details
251 can be found in Ref. [22].

252 D. Track finding with ACTS CKF

253 Starting from a set of initial track parameters, the ACTS
254 CKF is driven by the ACTS track propagator to search for
255 compatible measurements at a particular surface through KF
256 track fitting, as illustrated by Fig. 6. This process is also
257 known as track following. The measurement providing the
258 best fitting quality is associated to the track and used to filter
259 the track parameters for further track propagation.

260 IV. PERFORMANCE STUDIES

261 A. Monte-Carlo samples

262 The J/ψ decay process $J/\psi \rightarrow \Lambda(\rightarrow p\pi^-)\bar{\Lambda}(\rightarrow \bar{p}\pi^+)$ is
263 an important benchmark process at STCF allowing for several
264 important physics studies relevant to Λ . Those signal events
265 generated with the KKMC and EvtGen generators in OSCAR
266 without considering bream-related backgrounds are used to
267 evaluate the tracking performance. The 2D distributions of
268 the $\cos\theta$ versus p_T , vertex displacement in the x - y plane V_{xy}
269 versus p_T , and transverse impact track parameter d_0 versus
270 p_T , for proton (anti-proton), denoted as $p(\bar{p})$, and π in the
271 $J/\psi \rightarrow \Lambda(\rightarrow p\pi^-)\bar{\Lambda}(\rightarrow \bar{p}\pi^+)$ events are shown in Fig. 7.
272 The π has a low momentum with p_T below 310 MeV/c and

273 $p(\bar{p})$ has a p_T up to 1.1 GeV/c. A non-negligible amount of
274 particles are decaying outside the first layer of ITK. However,
275 most of the tracks have d_0 smaller than the radius of the first
276 ITK layer, in particular for $p(\bar{p})$.

277 Following event generation, Geant4 simulates hits from fi-
278 nal state particles decaying from primary particles interacting
279 with the STCF tracking system in a uniform magnetic field
280 of 1T. Detector measurements are then generated by apply-
281 ing Gaussian smearing to the positions of simulated hits, with
282 zero means and widths corresponding to the detector resolu-
283 tions.

284

B. Track finding performance

285 The performance of track finding, including seed finding
286 using either ACTS seeding algorithm or Hough Transform
287 algorithm at the first stage, and track following using ACTS
288 CKF at the second stage, is studied. Considering the accep-
289 tance of STCF tracking system, only truth particles with p_T
290 above 50 MeV and $|\cos\theta|$ below 0.94 are considered in the
291 performance metrics evaluation, which involves identifying
292 the primary particle [23] of a seed or a track, i.e. the simu-
293 lated particle which has the most simulated hits contributing
294 to this seed or track.

295 The seeding process serves as the initial step in track find-
296 ing using CKF, which should provide seeds for all particles
297 in an ideal case. The ACTS seeding efficiency is defined as
298 the fraction of particles in the tracking system acceptance re-
299 gion that have matched seeds with all three hits arising from
300 the same particle. The seeding efficiency using Hough Trans-
301 form is defined by requiring that a matched seed has at least
302 50% hits from its primary particle.

303 For ACTS seeding, it's only possible to find seeds for a
304 track if the particle produces hits in all three layers of ITK,
305 indicating a vertex displacement below 66.5 mm. The com-
306 parison between efficiencies of ACTS seeding and Hough
307 Transform algorithm as a function of V_{xy} of the particles is
308 shown in Fig. 8 top panels. The ACTS seeding efficiency ap-
309 proaches 100% when the number of measurements from ITK
310 is no less than 3. In particular, the ACTS seeding provides
311 better seeding efficiency than Hough Transform for π with
312 small V_{xy} . However, ACTS seeding efficiency immediately
313 drops to zero if the number of measurements from ITK is
314 below 3, indicating a significant limitation of ACTS seeding
315 algorithm, in particular for long-lived particles. The Hough
316 Transform algorithm, functioning as a global tracking algo-
317 rithm, demonstrates reduced sensitivity to the number of ITK
318 layers particles traverse. Fig. 8 bottom panels shows the seed-
319 ing efficiency as a function of particle p_T . It is observed that
320 Hough Transform algorithm can provide an efficiency above
321 90% for $p(\bar{p})$ with p_T above 350 MeV/c and above 80% for π
322 with p_T above 85 MeV/c, which is much improved compared
323 to ACTS seeding.

324 The reconstructed tracks are required to have at least five
325 measurements on the track and have reconstructed $|\cos\theta| <$
326 0.94. A reconstructed track is matched to its primary particle
327 if the fraction of its hits from its primary particle, i.e. track

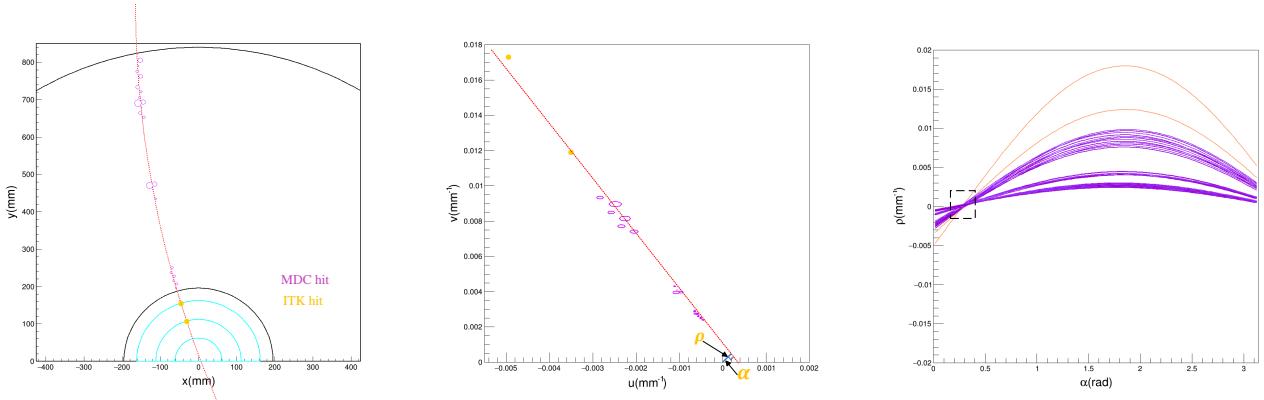


Fig. 4. An example of mapping detector measurements in geometrical transverse x - y plane (left) to the Conformal u - v space (middle) and eventually to the Hough curves in Hough space (right) for the particle p in a $J/\psi \rightarrow \Lambda(\rightarrow p\pi^-)\bar{\Lambda}(\rightarrow \bar{p}\pi^+)$ event. The particle decays between the first and the second layer of the ITK detector and hence has two hits on ITK. The hits or curves for ITK hits and MDC hits are shown by yellow and purple colors, respectively.

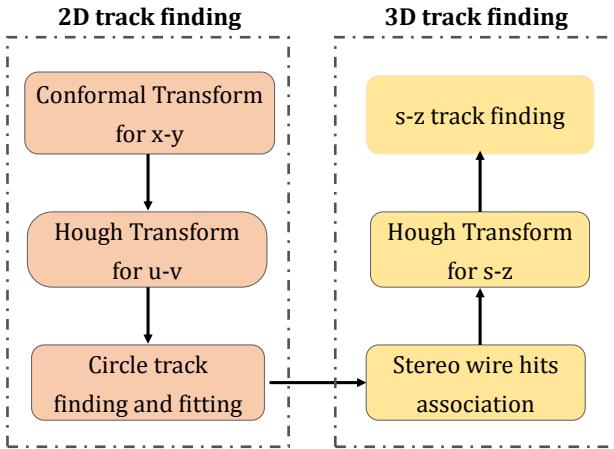


Fig. 5. The workflow of track finding using Hough Transform in OSCAR.

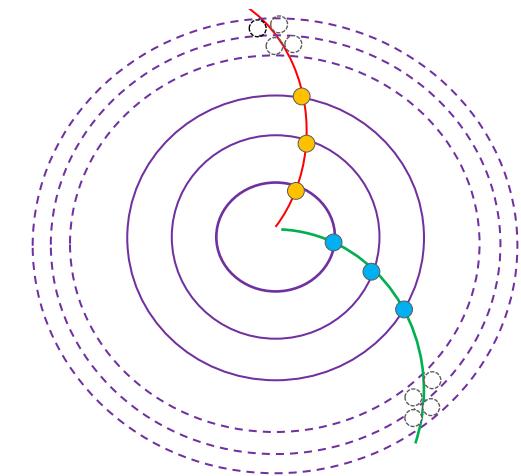


Fig. 6. Illustration of track finding using ACTS CKF with STCF ITK and MDC. Only two MDC layers are shown in the figure.

328 purity, is no less than 50%, and it's classified as a fake track
 329 if it's not matched to its primary particle. If more than one re-
 330 constructed tracks are matched to the same simulated particle,
 331 the track with the highest track purity is classified as the real
 332 track and others are classified as duplicate tracks. The track
 333 reconstruction efficiency is defined by the fraction of particles
 334 which have matched reconstructed tracks among the particles
 335 which have at least 5 simulated hits in the detector acceptance
 336 region. The fake rate is defined by the fraction of fake tracks
 337 among the reconstructed tracks. The duplicate rate is defined
 338 by the fraction of particles which have at least one duplicate
 339 track among the particles which have at least 5 simulated hits
 340 in the detector acceptance region.

341 Figure 9 shows the tracking efficiency as a function of par-
 342 ticle V_{xy} and p_T . As expected, the tracking efficiency us-
 343 ing ACTS seeding and CKF drops to zero when V_{xy} of the
 344 particle exceeds 66.5 mm, while the tracking efficiency with

345 Hough Transform and CKF is less dependent on the particle
 346 V_{xy} . A tracking efficiency above 80% for $p(\bar{p})$ with p_T above
 347 350 MeV/c and above 70% for π with p_T above 85 MeV/c is
 348 achieved using combined Hough Transform and CKF.

349 Figure 10 shows the fake rate and duplicate rate using the
 350 two different seeding strategies. The fake rate is less than
 351 0.4% and a non-negligible amount of duplicate tracks are
 352 found for particles with p_T below 150 MeV/c, which have
 353 looping trajectories when traversing the detector in a mag-
 354 netic field. ACTS seeding results in lower fake rate and du-
 355 plicate rate than that using Hough Transform as seeding.

356 The track finding performance of the alternative MAPS-
 357 based ITK for the long-lived particles is compared to that
 358 of the μ -RWELL-based ITK. Figure 11 shows the track-
 359 ing efficiency, fake rate and duplicate rate of the combined
 360 Hough Transform and ACTS CKF for the two different ITK
 361 designs. Since the first two layers of the MAPS-based

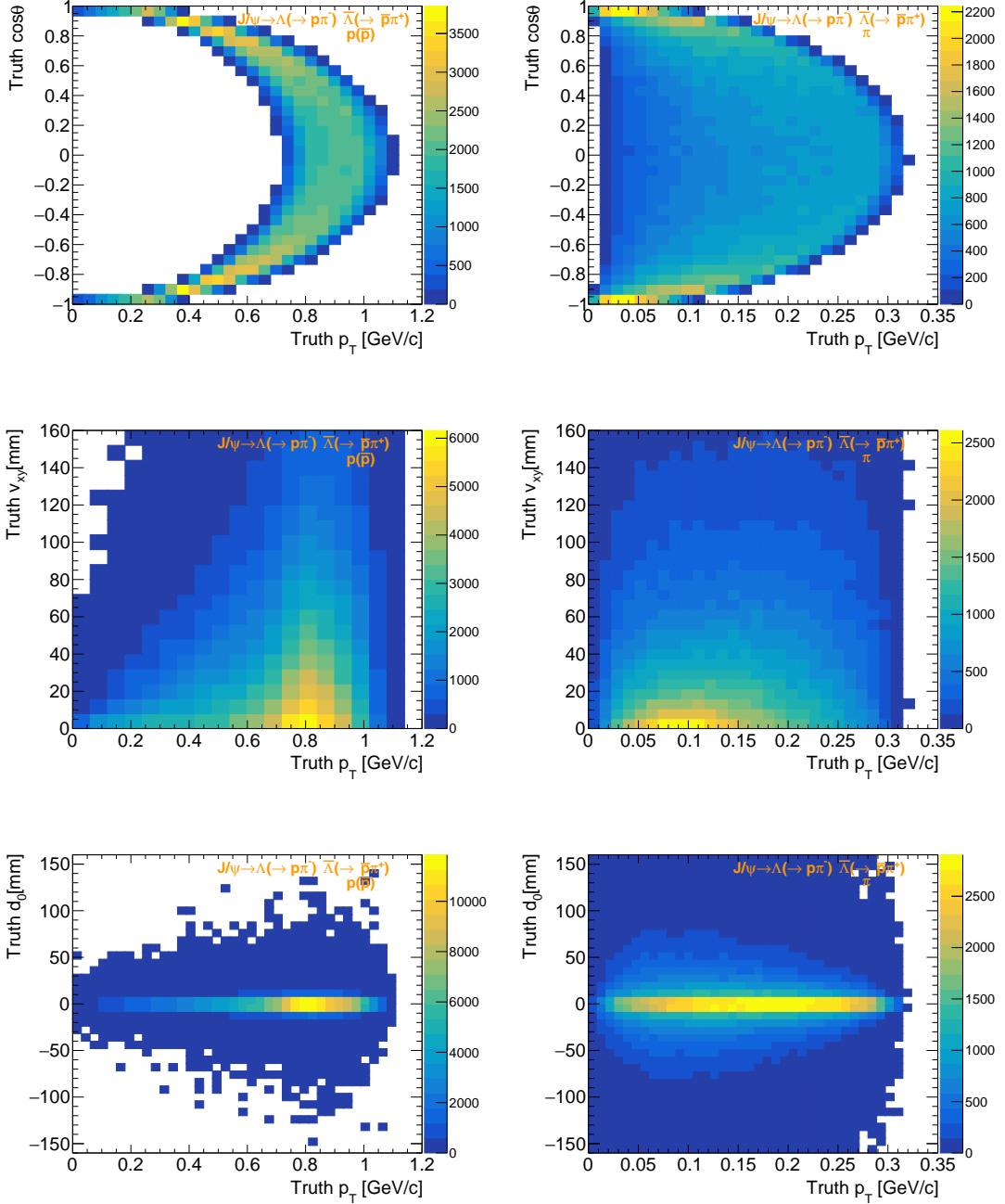


Fig. 7. The distributions of particle $\cos\theta$ versus p_T (top) and particle vertex displacement in the x - y plane V_{xy} versus p_T (middle), and d_0 versus p_T (bottom) for $p(\bar{p})$ (left) and π (right) in $J/\psi \rightarrow \Lambda(\rightarrow p\pi^-)\bar{\Lambda}(\rightarrow \bar{p}\pi^+)$ events.

362 ITK have smaller radii than the μ -RWELL-based ITK, the
 363 MAPS-based ITK is less robust against the long-lived parti-
 364 cles and provides slightly worse tracking efficiency than the
 365 μ -RWELL-based ITK. The occurrence of fake and duplicate
 366 tracks with the two designs is at similar level.

367

V. CONCLUSION

368 Processes with long-lived particles provide opportunities
 369 for probing CP, strong interaction etc. at the next generation
 370 of Tau-Charm facility, STCF. However, high-performance
 371 track reconstruction for long-lived particles is a challeng-
 372 ing and complicated task based on the tracking system of
 373 STCF. CKF is one of the most commonly used track find-

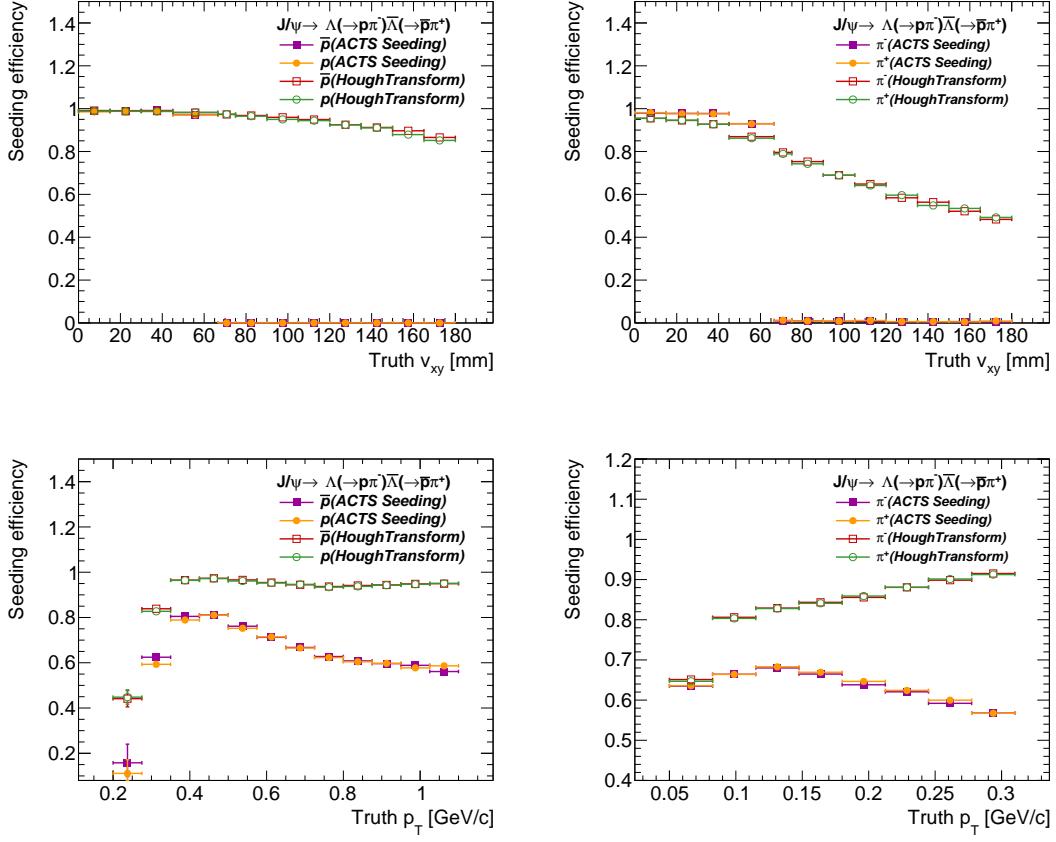


Fig. 8. The seeding efficiency as a function of the particle V_{xy} (top) and p_T (bottom) for $p(\bar{p})$ (left) and $\pi^+(\pi^-)$ (right) in 200k $J/\psi \rightarrow \Lambda(\rightarrow p\pi^-)\bar{\Lambda}(\rightarrow \bar{p}\pi^+)$ events. The solid purple square and yellow dot represent the results of ACTS seeding for particles with negative charge and positive charge, respectively. The hollow red square and green circle represent the results of Hough Transform for particles with negative charge and positive charge, respectively.

ing algorithms at HEP experiments with its performance subject to the performance of the corresponding seeding algorithm. For long-lived particles, CKF using traditional seeding strategy, which often uses measurements from inner detector(s), demonstrates significant performance loss. Based on the STCF offline software and the common tracking software ACTS, the combined performance of using Hough Transform as a seeding algorithm for ACTS CKF has been studied for the first time. The performance was evaluated using $J/\psi \rightarrow \Lambda(\rightarrow p\pi^-)\bar{\Lambda}(\rightarrow \bar{p}\pi^+)$ events at STCF. The study shows that CKF steered by Hough Transform ends up with improved efficiency compared to CKF steered by traditional seeding algorithm for particles with large vertex displacement. The tracking efficiency using combined Hough Transform and CKF is 80% for proton and anti-proton with p_T above 350 MeV/c, and above 70% for π with p_T above

85 MeV/c, with negligible occurrence of fake tracks. Duplicate tracks also exist, mainly arising from particles with p_T below 150 MeV/c with looping trajectories. Future development like extension of the 2D Hough space to 3D Hough space, where a track projection on the x - y plane not passing through origin is described by three dedicated parameters, is foreseen to further enhance the tracking efficiency for long-lived particles at STCF and beyond.

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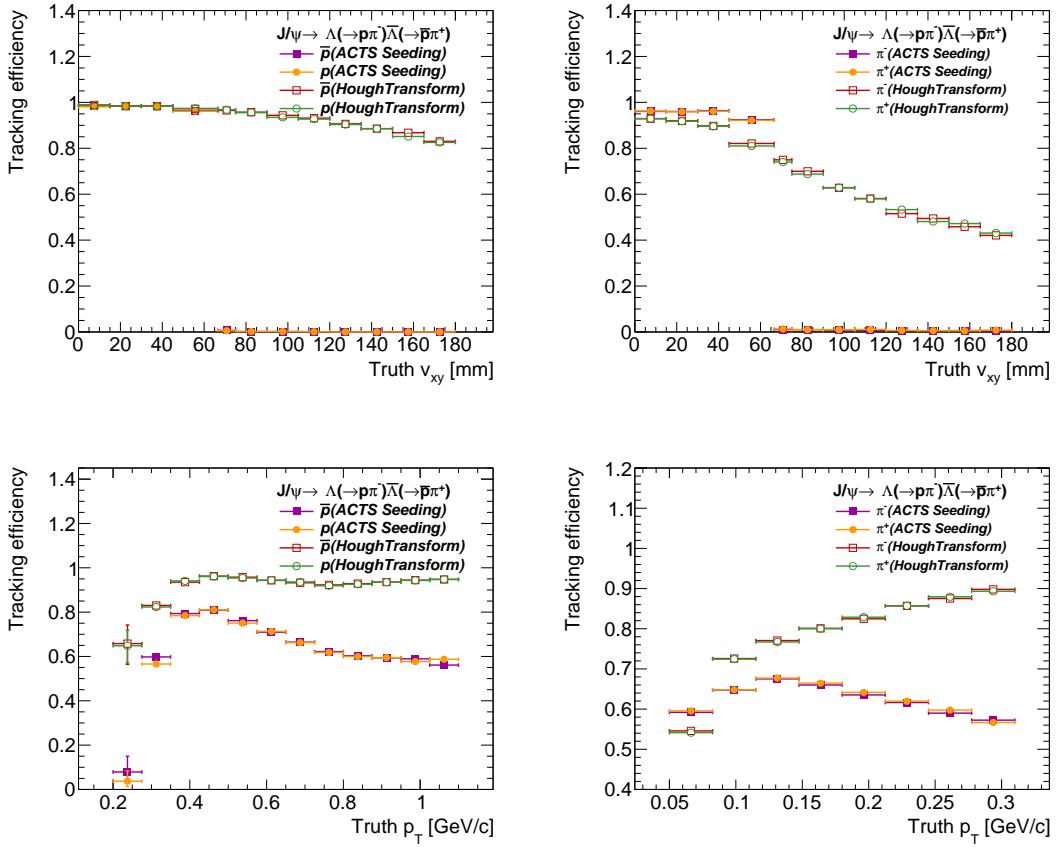


Fig. 9. The tracking efficiency as a function of the particle V_{xy} (top) and p_T (bottom) for $p(\bar{p})$ (left) and $\pi^+(\pi^-)$ (right) in 200k $J/\psi \rightarrow \Lambda(\rightarrow p\pi^-)\bar{\Lambda}(\rightarrow \bar{p}\pi^+)$ events. The solid purple square and yellow dot represent the results with ACTS seeding for particles with negative charge and positive charge, respectively. The hollow red square and green circle represent the results with Hough Transform for particles with negative charge and positive charge, respectively.

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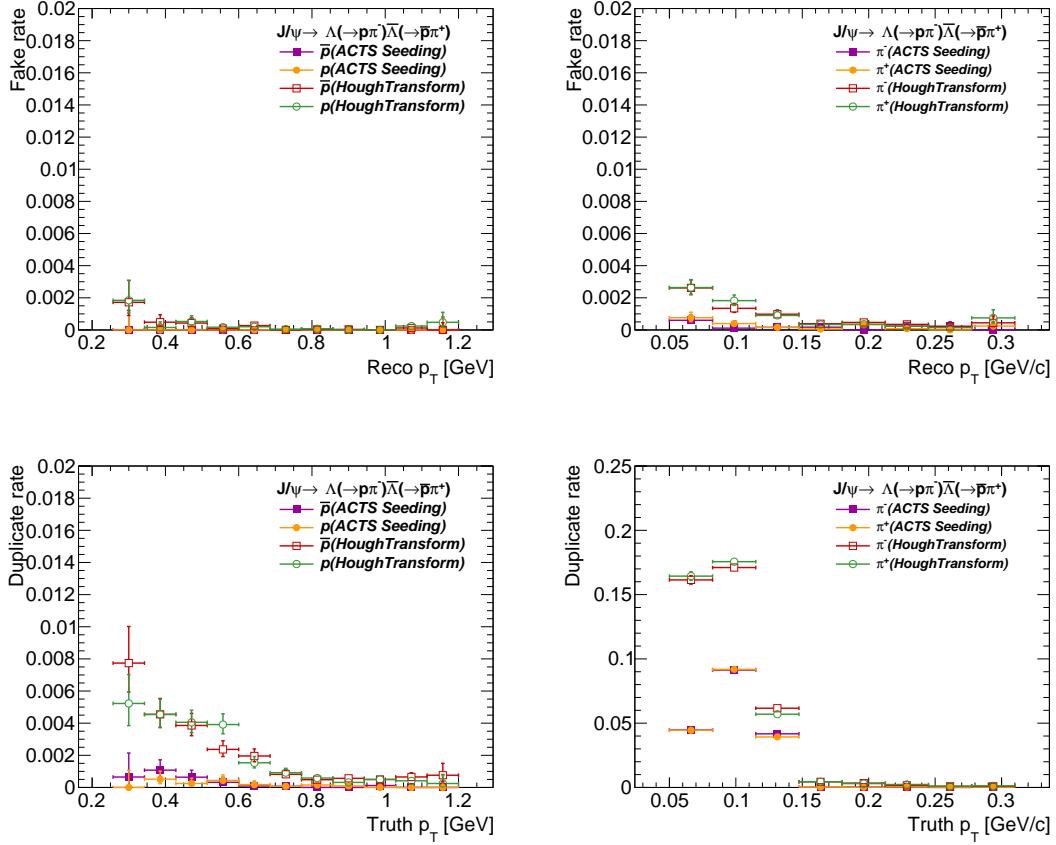


Fig. 10. The fake rate as a function of the track p_T (top) and duplicate rate as a function of the particle p_T (bottom) for $p(\bar{p})$ (left) and $\pi^+(\pi^-)$ (right) in 200k $J/\psi \rightarrow \Lambda(\rightarrow p\pi^-)\bar{\Lambda}(\rightarrow \bar{p}\pi^+)$ events. The solid purple square and yellow dot represent the results with ACTS seeding for particles with negative charge and positive charge, respectively. The hollow red square and green circle represent the results with Hough Transform for particles with negative charge and positive charge, respectively.

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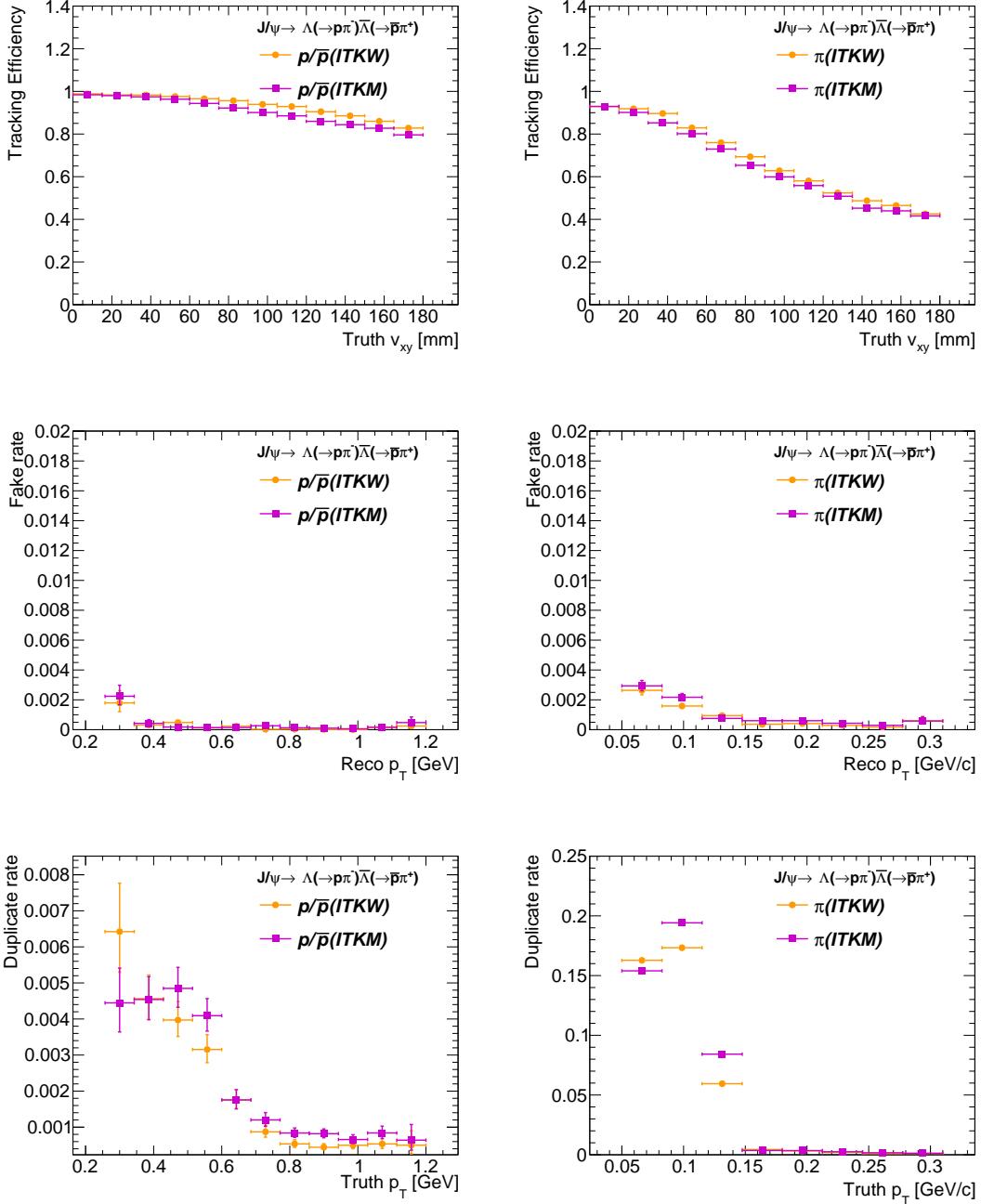


Fig. 11. The tracking efficiency as a function of the particle V_{xy} (top), fake rate as a function of the track p_T (middle) and duplicate rate as a function of the particle p_T (bottom) for $p(\bar{p})$ (left) and π (right) in 200k $J/\psi \rightarrow \Lambda(\rightarrow p\pi^-)\bar{\Lambda}(\rightarrow \bar{p}\pi^+)$ events using combined Hough Transform and ACTS CKF. The yellow dot and purple square represent the results with the μ -RWELL-based ITK (denoted as ITKW) and MAPS-based ITK (denoted as ITKM).

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